

$$\rho u \frac{\partial v}{\partial x} + \rho v \frac{\partial v}{\partial y} + \rho w \frac{\partial v}{\partial z} = \frac{\partial}{\partial y} \left(\lambda \nabla \cdot \mathbf{V} + 2\mu \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial z} \left[\mu \left(\frac{\partial w}{\partial y} + \frac{\partial v}{\partial z} \right) \right]$$

$$= E_1 + \sqrt{N} Q^{-1} (1 - P\alpha)$$

$$= \sqrt{N} [Q^{-1} (P\alpha) + Q^{-1} (1 - P\alpha)]$$

$$1 = \int_0^\infty \exp(-v^2/2N) dv$$

Van Hove singularities in intersubband transitions in multiquantum well photodetectors

**J. Le Rouzo^a, I. Ribet-Mohamed^a, R.Haidar, M. Tauvy^a, N. Guérineau^a,
E. Rosencher^{a,b} & S.L. Chuang^c**

^a ONERA / DOTA / CIO

^b Also Ecole Polytechnique, Palaiseau, France

^c University of Illinois

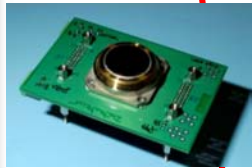
QWIP 2006 Workshop,
June 18-24, 2006 – Kandy, Sri Lanka

Context

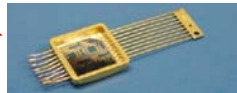
ONERA = French Agency for Aeronautics and Aerospace

Characterization of IR detectors

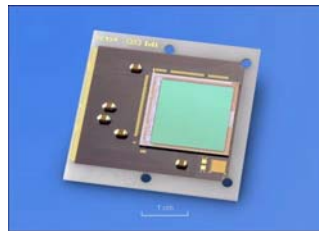
- Evaluation of the radiometric performances of IR sensors (MCT, Si:Ga, InSb, μ bolometers, QWIPs....)
- Realization of innovative test benches :
Angular response measurements,
MTF measurements,
Spectral characterization, ...
- Development of device models



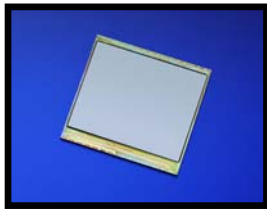
*μ bolometer
320x240 ULIS*



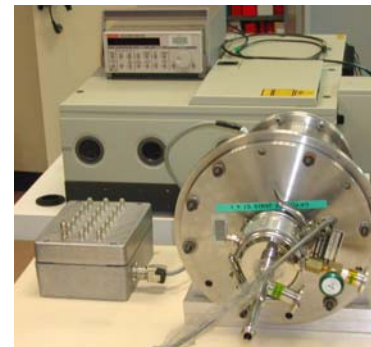
*Single element QWIP "low-noise"
Fraunhofer Institut*



*1000x1000 "IRCMOS" array of MCT
CEA/LIR*



*640x512 array of QWIP
Thales-R&T*



FTIR

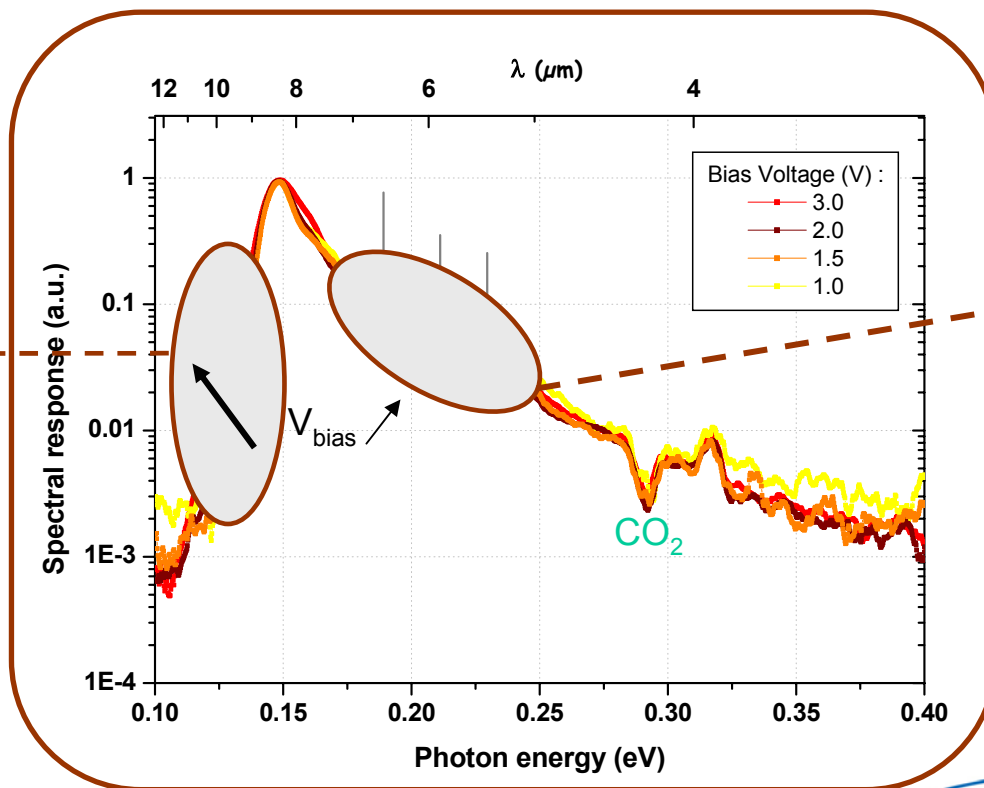


*"BRASIL"**

ONERA

Introduction

- Purpose : Develop new experiments in order to enhance our knowledge on QWIP Physics
- Idea : Log spectra of QWIP responsivity to allow off-band analysis



low energy part :
cut-off affected by
the applied electric
field

High energy part :
resonant features

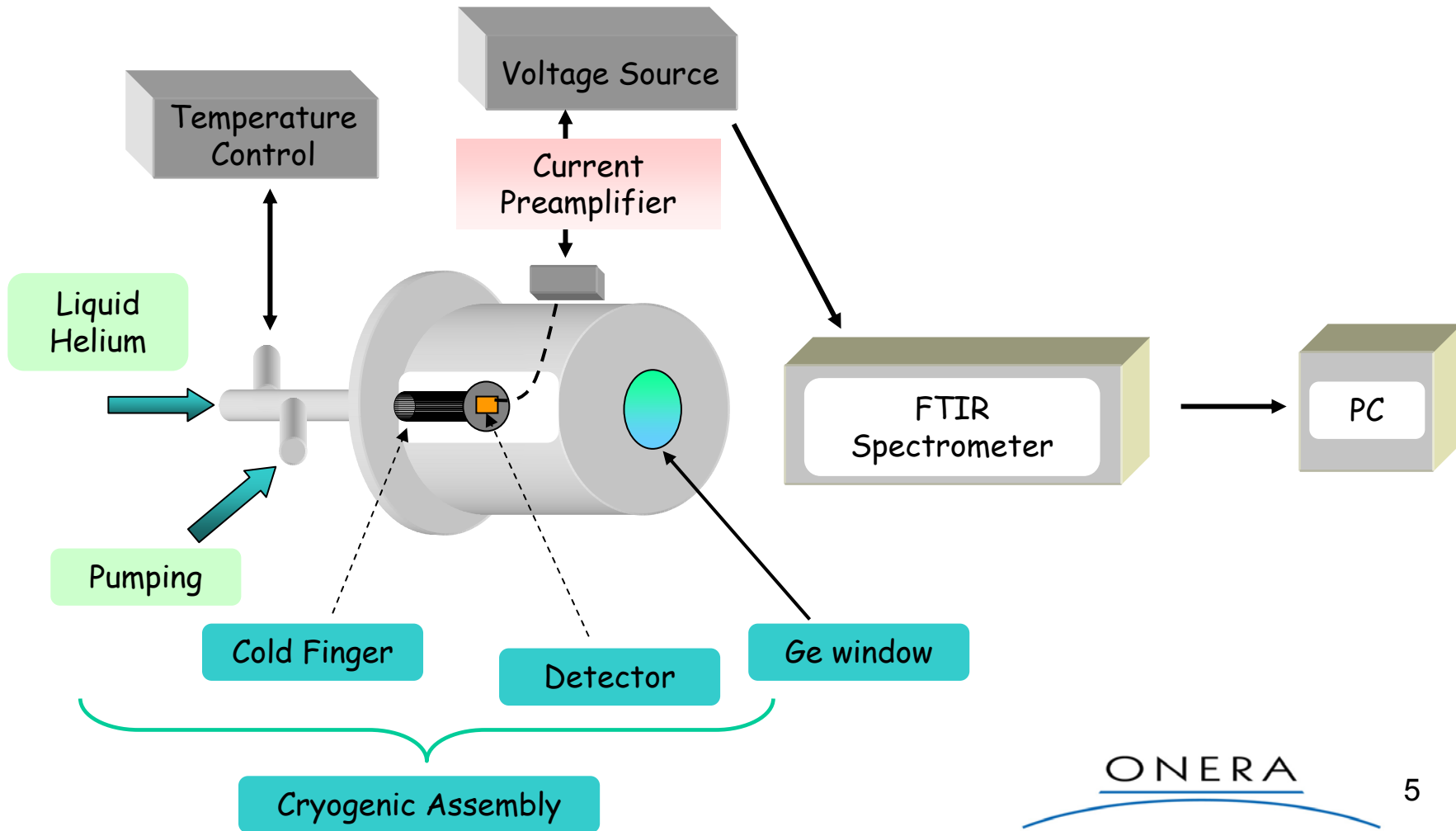
Outline

➡ Experimental setup

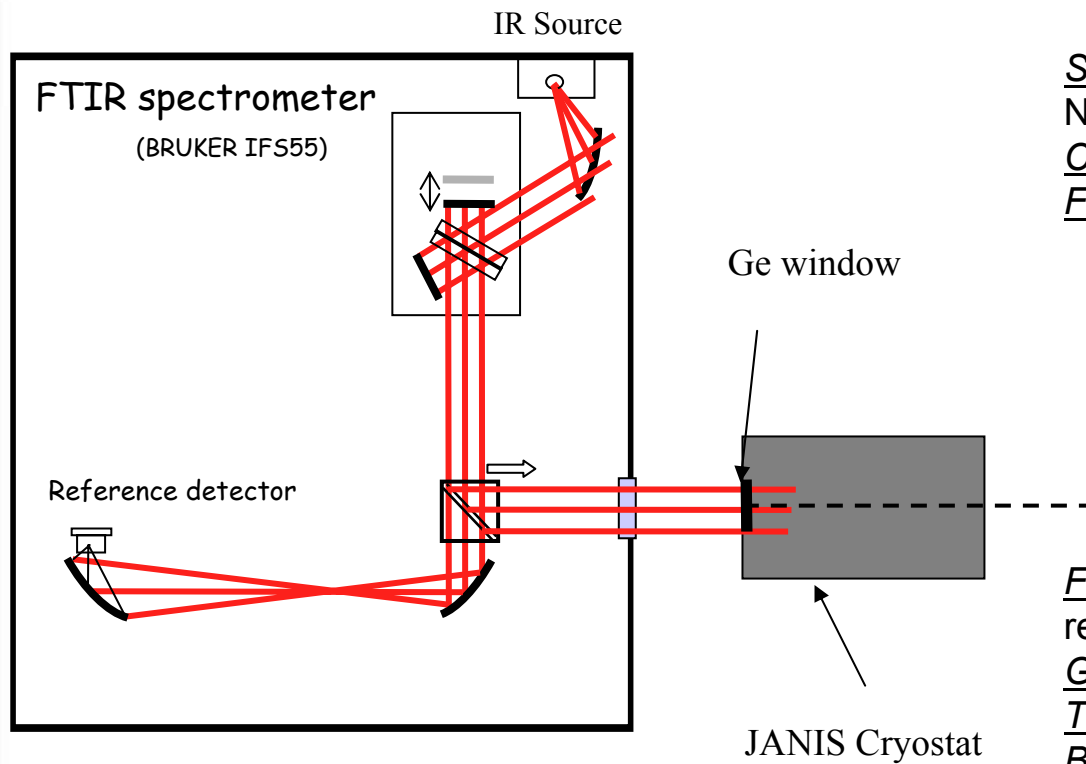
➡ Theoretical model

➡ Main results

• Test bench to perform spectral response measurements :



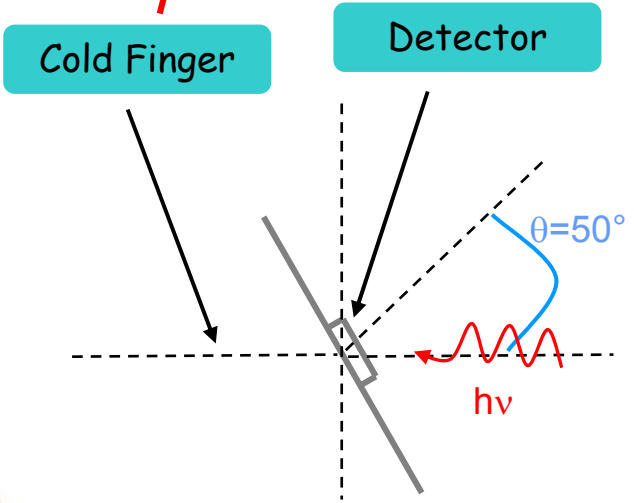
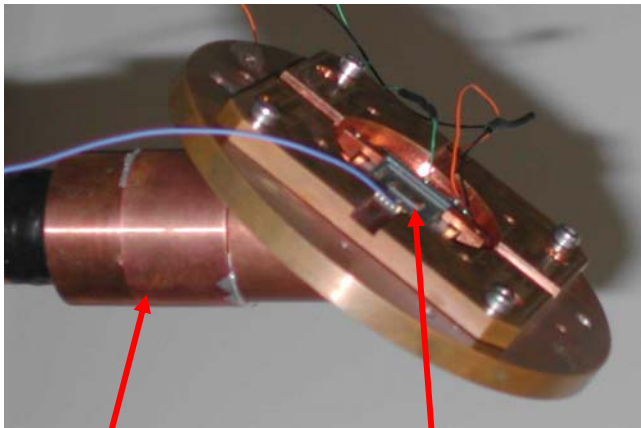
• FTIR spectrometer :



Sample : $\lambda=8.5 \mu\text{m}$,
 N=40 QWs, pixel without grating
Cryostat : JANIS
Fluid : liquid Helium

FTIR parameters : $f_{\text{acq}} = 5\text{kHz}$,
 resolution = 4cm^{-1}
Ge window ($\lambda_{\text{cut}} \sim 16\mu\text{m}$ (625cm^{-1}))
 $T_{\text{sample}} = 40\text{K}$
Bias voltage : -3, -2, -1.5, -1 V

Characteristics of the single detectors under study :

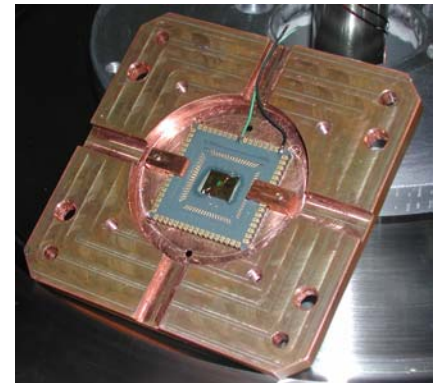


Sample :

QWIP Thales - RT
 $\lambda = 8.5 \mu\text{m}$,
 $\phi = 100 \mu\text{m}^2$
 $N = 40$ QWs
 &
pixel without grating

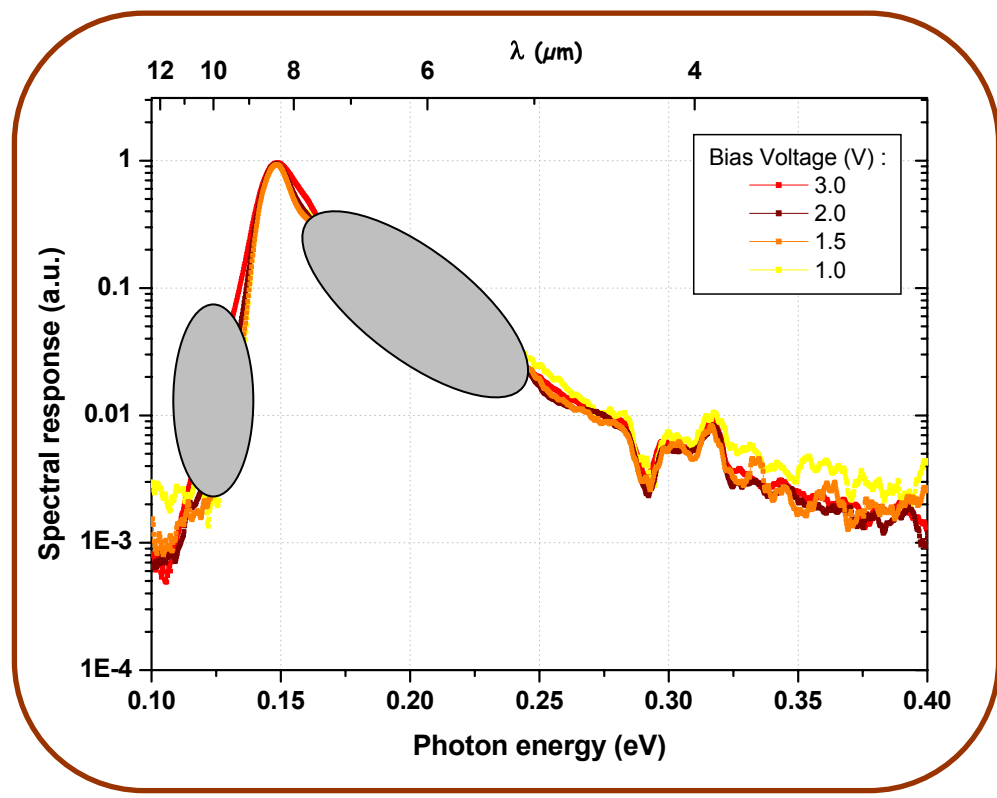


Pixel illuminate
 under
 tilted incidence



* Note : An incidence angle of 50° on the sample leads to an angle of 13.8° on the quantum wells

• Experimental results :



✓ Photocurrent spectra studied over more than 3 orders of magnitude

✓ Repeatability of the measurements

✓ Same spectral response shape for several angles of incidence on the sample

✓ New features in the spectral response at low and high energy !!!!

• Possible origin of high energy features ?

√ Calculation of QWIP absorption taking into account QW periodicity had never been done before

√ Solve numerically Schrödinger equation under applied electric field → Spurious resonances due to boundary conditions

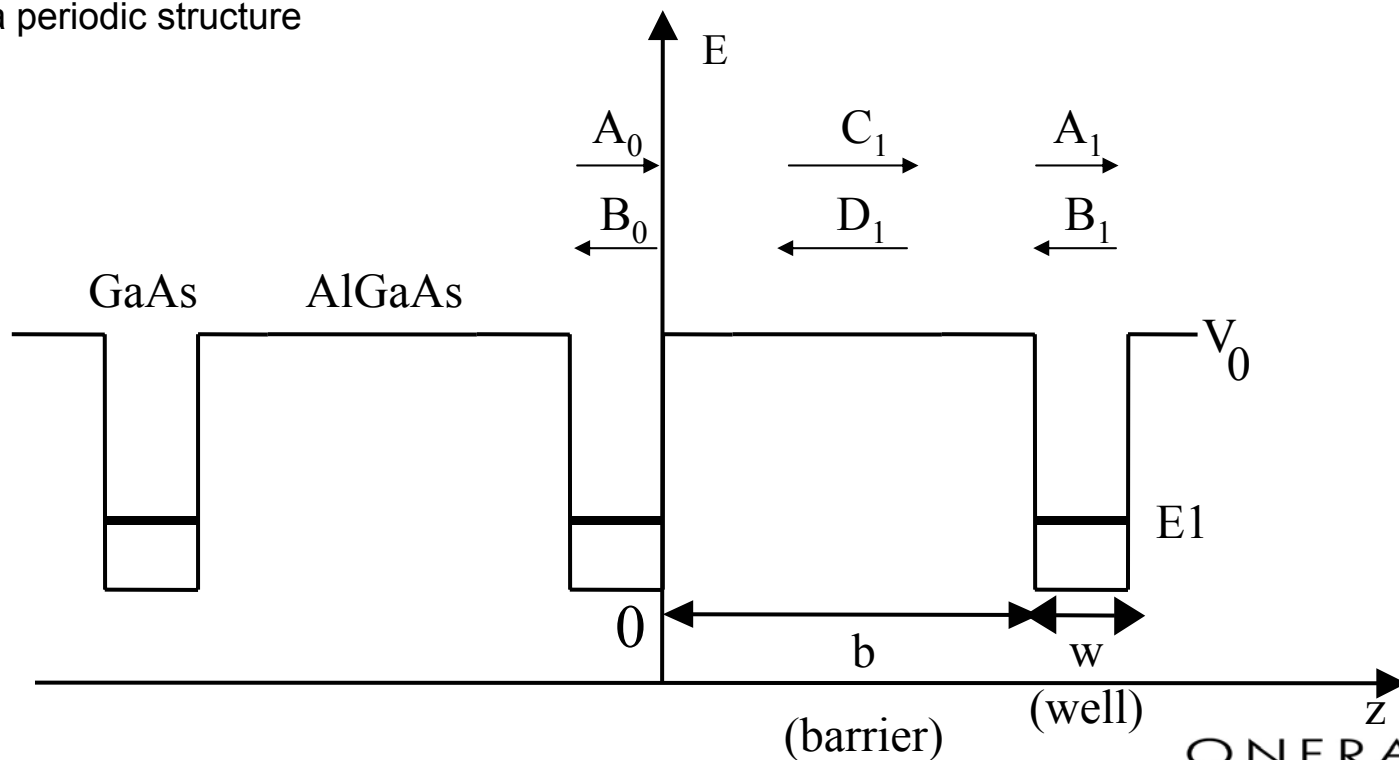
√ Features slightly affected by applied electric field

⇒ **Kronig-Penney
approach**

• Model :

1 Kronig-Penney model for multiquantum well structure :

√ Propagation matrix approach for a periodic structure



2 Schödinger equation for the envelope function*

$$\frac{-\hbar^2}{2m_{\text{eff}}}\Delta\Psi(z) + V\Psi(z) = E\Psi$$

$$\text{where } V = \begin{cases} V_0 & 0 \leq z < b \\ 0 & b \leq z < b + d = L \end{cases}$$

We get for each
period of the structure

$$\psi(z) = \begin{cases} A_0 e^{ikz} + B_0 e^{-ikz} & -d \leq z \leq 0 \\ C_1 e^{ik_b(z-d)} + D_1 e^{-ik_b(z-d)} & 0 \leq z \leq b \\ A_1 e^{ik(z-L)} + B_1 e^{-ik(z-L)} & b \leq z \leq L \end{cases}$$

Using the boundary conditions at $z=0$, we find the matrix equation for the coefficients of the wave functions in region 0 and the barrier b :

$$\begin{bmatrix} C_1 \\ D_1 \end{bmatrix} = \overline{\overline{F}}_{b0} \begin{bmatrix} A_0 \\ B_0 \end{bmatrix}$$

$$k = \sqrt{\frac{2m}{\hbar^2} E}$$

$$k_b = \sqrt{\frac{2m_b}{\hbar^2} (E - V_0)}$$

Similarly, the boundary conditions at $z=b$ give the matrix equation for region 1 and the barrier b :

$$\begin{bmatrix} A_1 \\ B_1 \end{bmatrix} = \overline{\overline{F}}_{1b} \begin{bmatrix} C_1 \\ D_1 \end{bmatrix}$$

We have the transition matrix for one period consisting of one well and one barrier :

$$\begin{bmatrix} A_1 \\ B_1 \end{bmatrix} = \overline{\overline{T}} \begin{bmatrix} A_0 \\ B_0 \end{bmatrix} \quad \text{where } \overline{\overline{T}} = \overline{\overline{F}}_{1b} \overline{\overline{F}}_{b0}$$

* S.L. Chuang, *Physics of Optoelectronic Devices* (Wiley Interscience, New York, 1995)

2 Schödinger equation for the envelope function

For the nth period, we find :

$$\begin{bmatrix} A_n \\ B_n \end{bmatrix} = \overline{T}^n \begin{bmatrix} A_0 \\ B_0 \end{bmatrix}$$

The eigenvalues and eigenvectors of the 2x2 matrix T are solutions of the determinantal equation :

$$\overline{T} \begin{bmatrix} A_0 \\ B_0 \end{bmatrix} = t \begin{bmatrix} A_0 \\ B_0 \end{bmatrix}$$

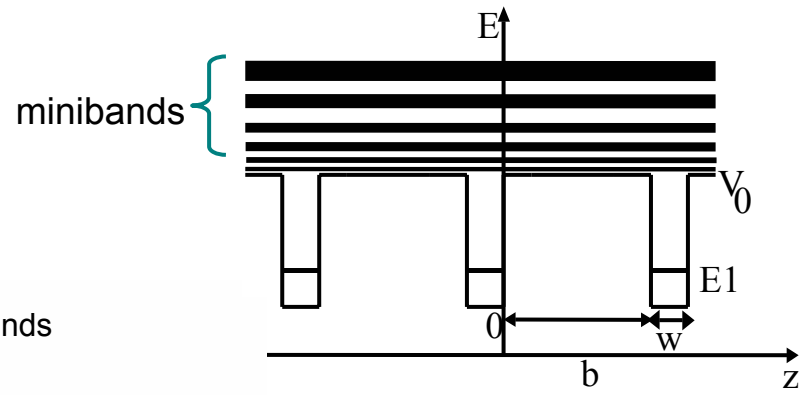
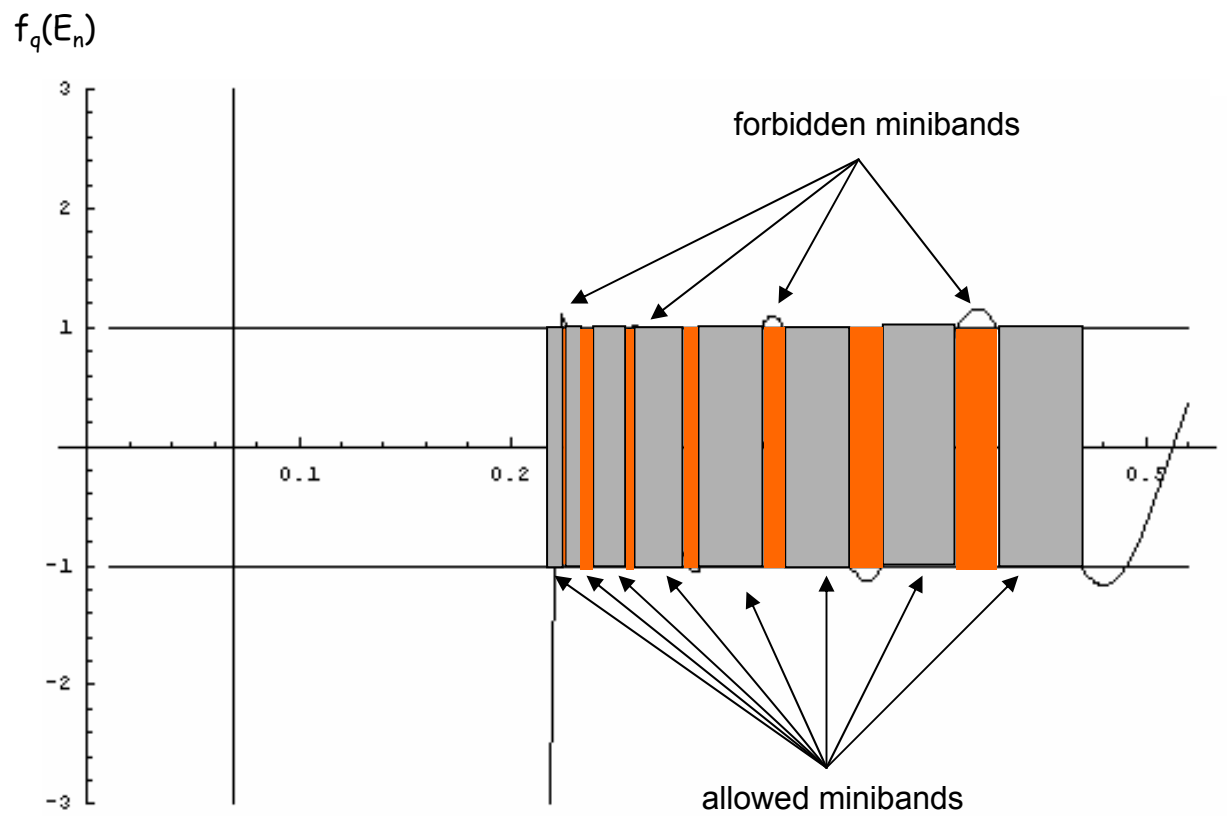
$$P = \frac{mk_b}{m_b k}$$

We get that the eigenequation is :

$$\cos(qL) = \cos(kw) \cos(k_b b) - \frac{1}{2} \left(\text{●} + \frac{1}{P} \right) \sin(kw) \sin(k_b b) = f_q(E)$$

3 Miniband :

→ condition on the energy levels : $|f_q(E_n)| \leq 1$



energy (eV)

4 Miniband dispersion :

- Bloch quantum state :

$$|n, k\rangle = u_{n,k}(z)e^{ikz}$$

Displays the periodicity of the QWIP structure

- Momentum matrix element between the bound state and the excited state:

$$p_{1n}(k) = \left\langle 1, k \left| \frac{\hbar}{i} \frac{d}{dz} \right| n, k \right\rangle \Rightarrow p_{bn}(k) = \hbar \int u_{1,k}(z)^* \frac{d}{dz} u_{n,k} dz$$

- Fermi Golden Rule :

$$P(h\nu) \propto \rho(E1 + h\nu) p_{b,n}(k)^2$$

- Absorption :

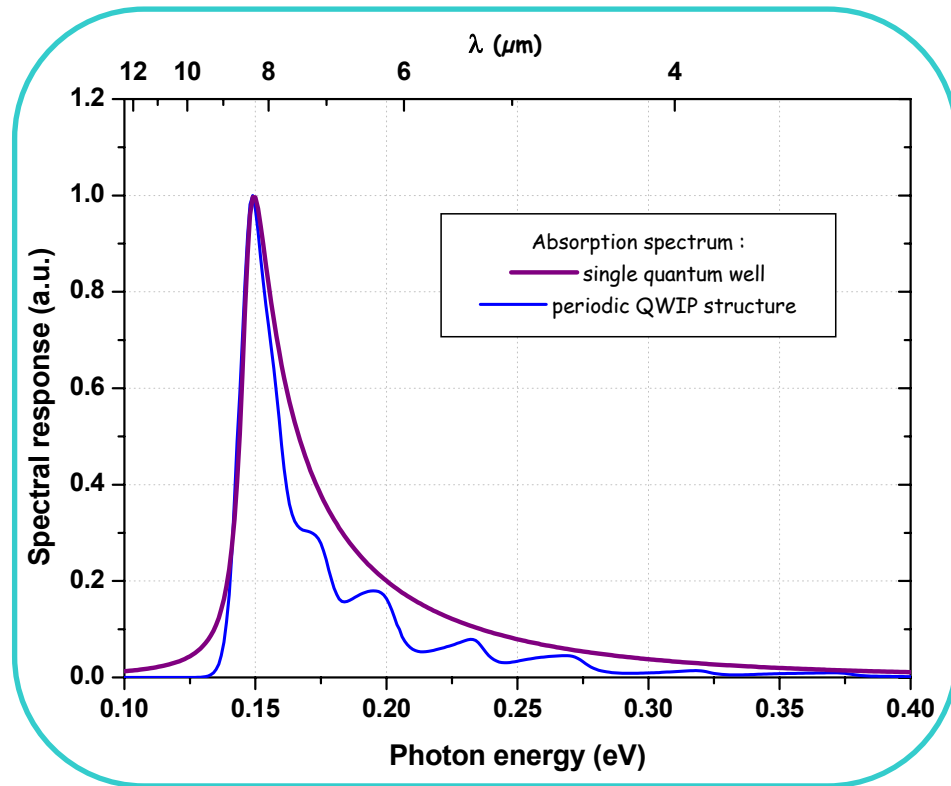
$$\alpha \propto P(h\nu) * \text{Lorentzian}(\Gamma, h\nu)$$

- Model : overview

Remind the approach of our study

- . Periodicity of the QWIP structure
- . Kronig-Penney approach
- . Response modeled by taking into account quantum absorption
- . Not included the already well known transport mechanisms
- . Empirical parameters needed like x_{Al} , b and w , Γ

• Importance of QW periodicity :



Purple line :

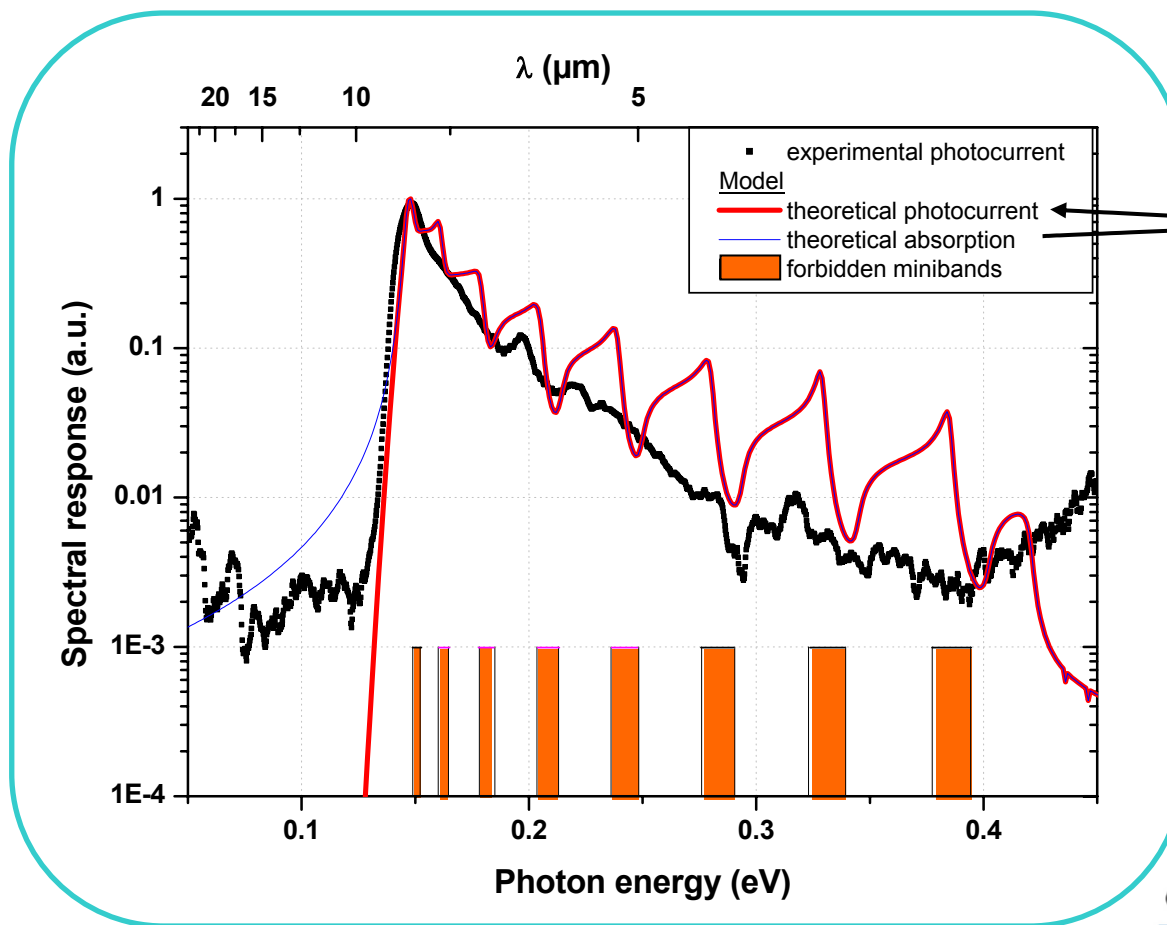
absorption spectrum obtained using the oscillator strength for the bound to continuum transition*

Blue line :

our Kronig-Penney approach

* H.C. Liu, J. Appl. Phys., 73, pp. 1-7 (1993)

- Comparison between an experimental QWIP log spectrum and our theoretical model :



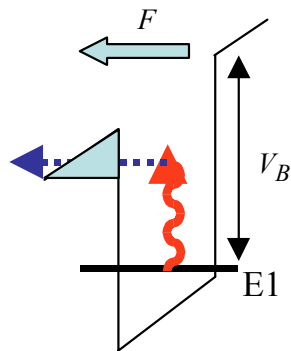
Taking into account of electric field assisted tunnelling

We noticed that the off-band response is non negligible (~1% in the 3–5 μm for a 8–12 μm detector)

• Two important results *

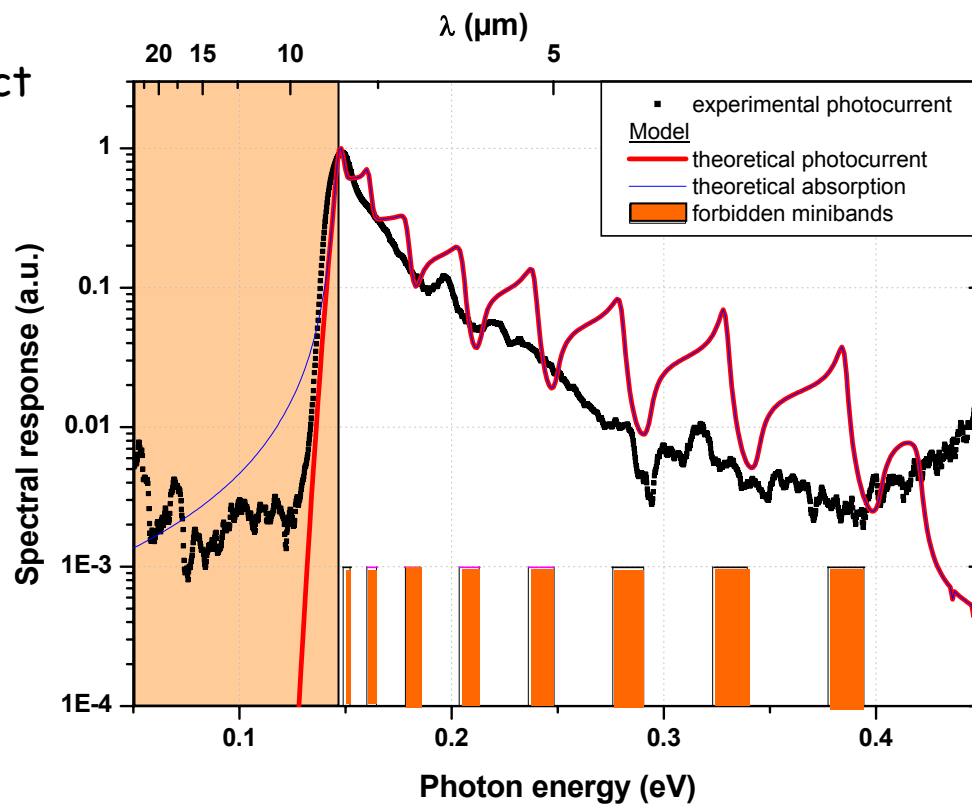
1

Low energy part :
Fowler-Nordheim effect



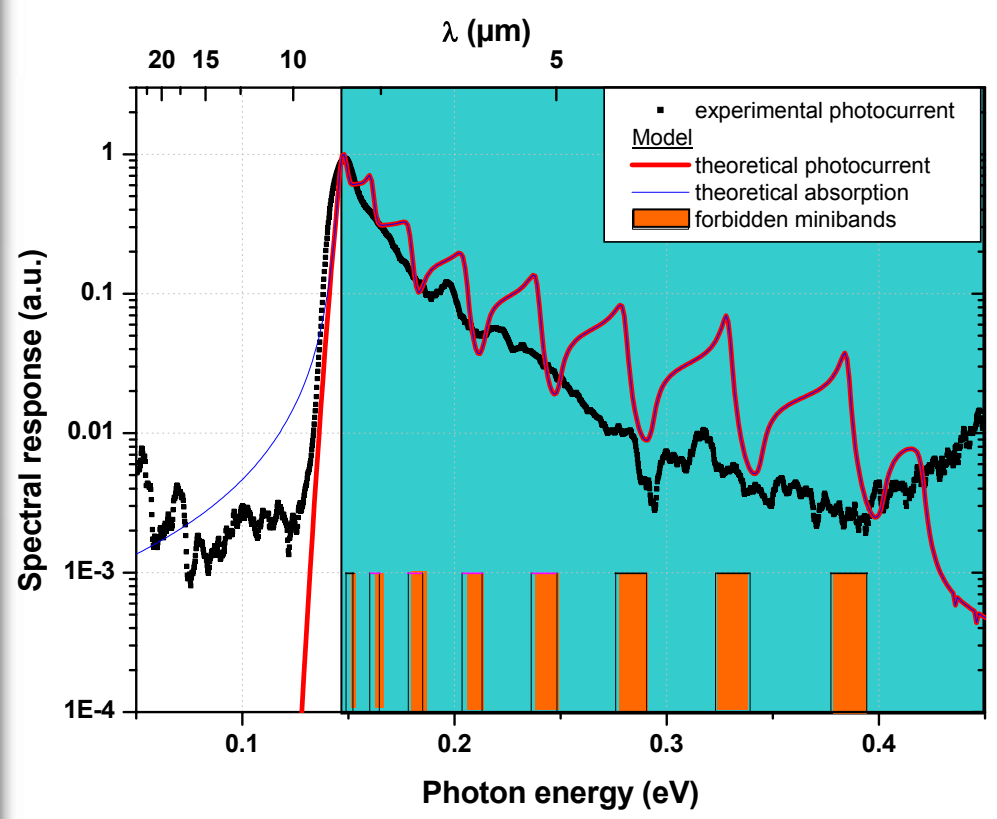
$V_B - E1$: Ionization threshold

➔ Electric field assisted tunnelling

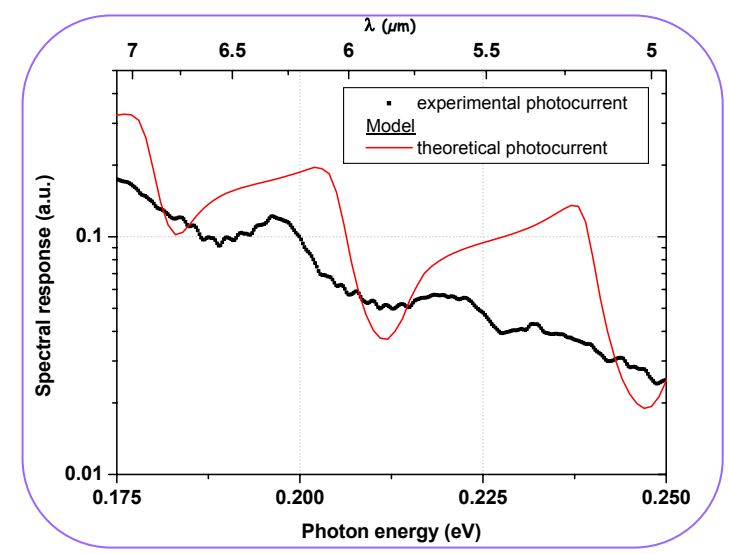


* J. Le Rouzo & al., *Appl. Phys.Lett*, **88**, 091117 (2006)

2



Van Hove singularities at the miniband edges



➔ Critical points in the density of states

Conclusion

**Electric field assisted
tunnelling**

&

Van Hove singularities

**Phenomena which have been
largely overlooked before**

**→ important role in the low
and high energy parts of the
QWIP photocurrent!**

Authors are also grateful to colleagues from
Thales-R&T for providing the samples

Thank you for your attention ...
Any question ?