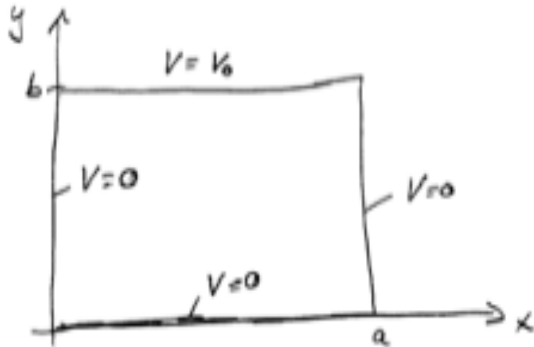




Solutions for Homework # 7

Problem #1: (30 points) An infinitely long rectangular conductive cylinder has three sides grounded, while the fourth (imagined to be insulated from others by very fine gaps) is held at potential V_0 . Find the potential at all interior points.!



$$\begin{aligned} x=0, & V=0 \\ x=a, & V=0 \\ y=0, & V=0 \\ y=b, & V=V_0 \end{aligned}$$

Laplace Equ. $\Delta V = 0$

Two dim. problem $V = V_x(x) \cdot V_y(y)$

$$\Rightarrow \Delta V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} = \frac{\partial^2 V_x(x)}{\partial x^2} \cdot V_y + \frac{\partial^2 V_y}{\partial y^2} \cdot V_x(x) = 0$$

$$\Rightarrow \Delta V_x + \kappa^2 V_x = 0 \quad \text{and} \quad \Delta V_y - \kappa^2 V_y = 0$$

$$\hookrightarrow V_x = A \cdot \sin(\alpha x) + B \cdot \cos(\alpha x)$$

$$V_y = C \cdot e^{\kappa y} + D \cdot e^{-\kappa y}$$

Boundary condition

$$V_x(x=0) = 0 = A \cdot \sin(0) + B \cdot \cos(0) \Rightarrow B = 0$$

$$V_x(x=a) = 0 = A \cdot \sin(\alpha a) \Rightarrow \alpha \cdot a = n \cdot \pi \Rightarrow \alpha = \frac{n\pi}{a}$$

$$\hookrightarrow V_x(x) = A_n \cdot \sin\left(\frac{n\pi}{a} \cdot x\right)$$

$$V_y(y=0) = 0 = C \cdot e^0 + D \cdot e^{-0} \Rightarrow C = -D$$

$$V_y(y=b) = V_0 = C \cdot [e^{\kappa b} - e^{-\kappa b}] = C \cdot \sinh(\kappa \cdot b)$$

$$\hookrightarrow \text{general solution} \quad V = V_x \cdot V_y = \sum_{n=1}^{\infty} \left[A_n \cdot \sin\left(\frac{n\pi}{a} \cdot x\right) \right] \cdot \left[C_n \cdot \sinh\left(\frac{n\pi}{a} \cdot y\right) \right]$$

\rightarrow

back to boundary conditions:

$$V(x, y=b) = V_0 = \sum_{n=1}^{\infty} \underbrace{A_n \cdot C_n}_{E_n} \cdot \sin\left(\frac{n\pi}{a}x\right) \cdot \sinh\left(\frac{n\pi}{a}b\right) \quad \Bigg/ \quad \int_0^a \sin\left(\frac{n'\pi}{a}x\right) dx$$

$$\rightarrow V_0 \int_0^a \sin\left(\frac{n'\pi}{a}x\right) dx = \sum_{n=1}^{\infty} E_n \cdot \sinh\left(\frac{n\pi}{a}b\right) \cdot \underbrace{\int_0^a \sin\left(\frac{n\pi}{a}x\right) \sin\left(\frac{n'\pi}{a}x\right) dx}_{= \frac{a}{2} \delta(n-n')}$$

$$\begin{aligned} \hookrightarrow V_0 \int_0^a \sin\left(\frac{n'\pi}{a}x\right) dx &= E_n' \sinh\left(\frac{n'\pi}{a}b\right) \cdot \frac{a}{2} \\ &= \frac{\cos\left(\frac{n'\pi}{a}x\right)}{\left(\frac{n'\pi}{a}\right)} \Bigg|_0^a = E_n' \frac{a}{2} \sinh\left(\frac{n'\pi b}{a}\right) \end{aligned}$$

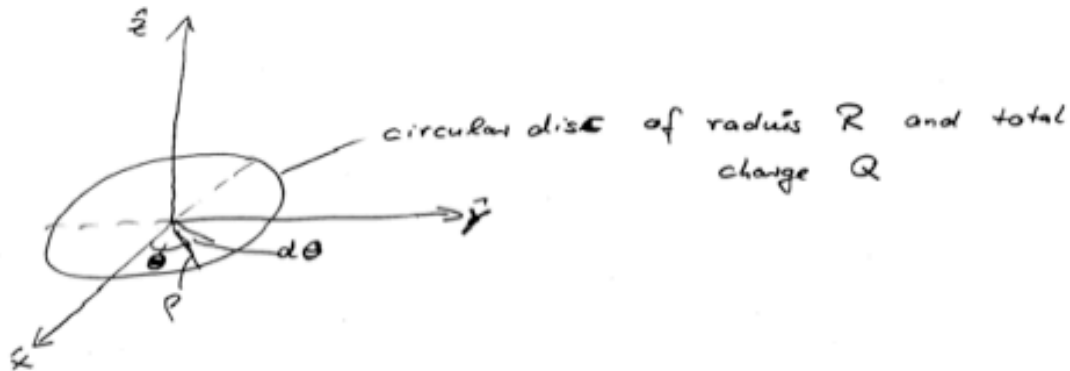
For $n' = \text{even integer}$: $E_n' = 0$

$n' = \text{odd}$ " : $E_n' = \frac{4V_0}{n'\pi} \left[\sinh\left(\frac{n'\pi b}{a}\right) \right]$

\hookrightarrow compl. solution

$$\phi(x, y) = \sum_{n \text{ odd}} \frac{4V_0}{n\pi} \frac{1}{\sinh\left(\frac{n\pi b}{a}\right)} \cdot \sin\left(\frac{n\pi}{a}x\right) \cdot \sinh\left(\frac{n\pi}{a}y\right)$$

Problem #2: (20 points): Find the potential and the electric field strength along the axis of a thin uniformly charged circular disc of radius R and total charge q . Show that the normal component of the field changes by σ/ϵ_0 on passing through the surface of the disc. Consider the field at large distances from the disc.



$$\text{potential } V(r) = \frac{1}{4\pi\epsilon_0} \frac{\sigma(r \cdot d\rho \cdot d\theta)}{|\vec{r}|}$$

$$\begin{aligned} V(z) &= \int_0^{2\pi} \int_0^R V(r) = \int_0^{2\pi} \int_0^R \frac{1}{4\pi\epsilon_0} \frac{\sigma \rho \cdot d\rho \cdot d\theta}{\sqrt{\rho^2 + z^2}} \\ &= \frac{\sigma}{4\pi\epsilon_0} 2\pi \sqrt{\rho^2 + z^2} \Big|_0^R = \frac{\sigma}{2\epsilon_0} [\sqrt{R^2 + z^2} - z] \end{aligned}$$

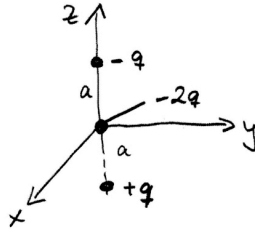
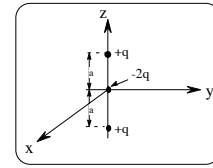
$$\begin{aligned} E = -\nabla V \Rightarrow E_z &= \frac{\partial}{\partial z} V \cdot \hat{z} = -\frac{\sigma}{2\epsilon_0} \left[\frac{1}{2} (R^2 + z^2)^{-1/2} \cdot 2z - 1 \right] \hat{z} \\ &= \frac{\sigma}{2\epsilon_0} \left[1 - \frac{z}{\sqrt{R^2 + z^2}} \right] \cdot \hat{z} \end{aligned}$$

$$E_{0+} = \frac{\sigma}{2\epsilon_0} \cdot \hat{z}, \quad E_{0-} = -\frac{\sigma}{2\epsilon_0} \hat{z}$$

$$\hookrightarrow |E_{0+} - E_{0-}| = \frac{\sigma}{2\epsilon_0} - \left(-\frac{\sigma}{2\epsilon_0}\right) = \underline{\underline{\frac{\sigma}{\epsilon_0}}}$$

Problem #3: (20 points)

- (a) Find the first two sets of non-vanishing moments, q_{lm} , for the following charge distribution
- (b) Write down the multipole expansion of the potential due to this charge distribution, keeping only the first two sets of q_{lm} .



$$\rho(\vec{r}) = \frac{q}{2\pi a^2} \left[-\delta(r) + \delta(r-a) \cdot \delta(\cos\theta-1) + \delta(r-a) \delta(\cos\theta+1) \right]$$

$$q_{lm} = \int d^3r' \rho(r') r'^l Y_{lm}^*(\theta', \varphi')$$

a.) We have azimuthal symmetry : $m=0$

$$Y_{lm}(\theta, 0) = \sqrt{\frac{2l+1}{4\pi}} \cdot P_l(\cos\theta)$$

$$\hookrightarrow q_{l0} = \int d^3r' \rho(r') r'^l \sqrt{\frac{2l+1}{4\pi}} P_l(\cos\theta)$$

$$\text{for } l=0: q_{00} = \frac{1}{\sqrt{4\pi}} \cdot \frac{q}{2\pi} \int \frac{1}{r'^2} \left[-\delta(r') + \delta(r'-a) \delta(\cos\theta'-1) + \delta(r'-a) \delta(\cos\theta'+1) \right] \cdot P_0(\cos\theta') \cdot r'^2 \cdot dr' \cdot \sin\theta' d\theta' d\varphi'$$

$$= \frac{q}{4\pi} \left[\int_0^a \delta(r') dr' \int_0^\pi \sin\theta' d\theta' + \int_0^a \delta(r'-a) dr' \int_0^\pi \delta(\cos\theta'-1) \sin\theta' d\theta' + \int_0^a \delta(r'-a) dr' \int_0^\pi \delta(\cos\theta'+1) \sin\theta' d\theta' \right]$$

$$= -2 + 1 + 1 = 0 : \text{monopole (net charge = zero)}$$

non-vanishing moments may be found from the general term for

$$q_{l0}, l \neq 0$$

$$q_{l0} = \frac{q}{2\pi} \sqrt{\frac{2l+1}{4\pi}} \int_0^{2\pi} d\varphi \int_0^\pi \int_0^a \left[-\delta(r') + \delta(r'-a) \delta(\cos\theta'-1) + \delta(r'-a) \delta(\cos\theta'+1) \right] \times P_l(\cos\theta') r'^l dr' \sin\theta' d\theta' \rightarrow$$

$$\rightarrow q_{l0} = \sqrt{\frac{2l+1}{4\pi}} \cdot q \cdot \left[(0)^l + a^l P_l(1) + a^l P_l(-1) \right]$$

$$= q \sqrt{\frac{2l+1}{4\pi}} a^l \left[\underbrace{P_l(1) + P_l(-1)}_{\substack{= 0 \text{ for odd } l \\ = 2 \text{ for even } l}} \right]$$

$$= 2q \sqrt{\frac{2l+1}{4\pi}} a^l \text{ for even } l \quad (\text{since } P_l(1) = 1 \text{ for all } l)$$

$$l=2: q_{20} = 2q \sqrt{\frac{5}{4\pi}} a^2 = q a^2 \cdot \sqrt{\frac{5}{\pi}} \quad (*) \rightarrow$$

$$l=4: q_{40} = 2q \sqrt{\frac{9}{4\pi}} a^4 = q \cdot a^4 \cdot \frac{3}{\sqrt{\pi}}$$

$$b.) \quad V(\vec{r}) = 4\pi \sum_l \frac{1}{2l+1} q_{l0} \frac{1}{r^{l+1}} \sqrt{\frac{2l+1}{4\pi}} \cdot P_l(\cos\theta) \quad | \quad m=0!$$

$$l=2,4: \rightarrow V(\vec{r}) = 4\pi \left[\frac{1}{5} q_{20} \frac{1}{r^3} \sqrt{\frac{5}{4\pi}} P_2(\cos\theta) + \frac{1}{9} q_{40} \frac{1}{r^5} \sqrt{\frac{9}{4\pi}} P_4(\cos\theta) + \dots \right]$$

$$= 4\pi \cdot \left[\frac{1}{5} q a^2 \sqrt{\frac{5}{\pi}} \cdot \sqrt{\frac{5}{4\pi}} P_2(\cos\theta) + \frac{1}{9} q \cdot a^4 \frac{3}{\sqrt{\pi}} \frac{1}{r^5} \sqrt{\frac{9}{4\pi}} P_4(\cos\theta) + \dots \right]$$

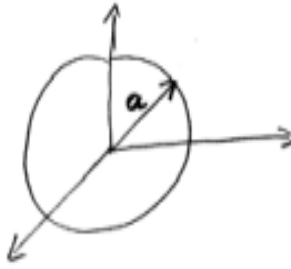
$$\rightarrow V(\vec{r}) = \frac{2q}{a} \left[\left(\frac{a}{r}\right)^3 P_2(\cos\theta) + \left(\frac{a}{r}\right)^5 P_4(\cos\theta) + \dots \right]$$

⊛ Momentums:

$$P = \int \vec{r}' \cdot \rho(\vec{r}') d^3r' = \dots = 0$$

$$Q_{ij} = \int \rho(\vec{r}') (3x'_i x'_j - \delta_{ij} r'^2) d^3r' = \dots \begin{pmatrix} -qa^2 & & 0 \\ & -qa^2 & \\ 0 & & 2qa^2 \end{pmatrix}$$

Problem #4: (20 points) - Jackson problem 3.5. A hollow sphere of an inner radius 'a' has the potential specified on its surface to be $V_0(\theta, \varphi)$. Prove the equivalence of the two forms of the solution for the potential inside the sphere



a) $G(\vec{r}, \vec{r}') = \frac{1}{|\vec{r} + \vec{r}'|^2 - 2\vec{r}\vec{r}'\cos\gamma} = \frac{1}{\frac{r^2+r'^2}{a^2} + a^2 - 2\vec{r}\vec{r}'\cos\gamma}$

$$\frac{\partial G}{\partial n'} = -\frac{r^2 - a^2}{a(r^2 + a^2 - 2\vec{r}\vec{r}'\cos\gamma)^{3/2}}$$

$$\text{Potential } V(\vec{r}) = -\frac{1}{4\pi} \oint V(\vec{r}') \frac{\partial G}{\partial n'} d^2r' = \frac{1}{4\pi} \int \frac{V(\theta', \varphi') a (a^2 - r^2)}{(r^2 + a^2 - 2ar\cos\gamma)^{3/2}} d\Omega'$$

$$= \frac{a(a^2 - r^2)}{4\pi} \int \frac{V(\theta', \varphi')}{(r^2 + a^2 - 2ar\cos\gamma)^{3/2}} d\Omega'$$

b.) we know that

$$V(r, \theta, \varphi) = \sum_{l,m} [A_{lm} r^l + B_{lm} r^{-(l+1)}] \cdot Y_{lm}(\theta, \varphi)$$

$$\text{at } r=a : V = V_0(\theta, \varphi)$$

$$\text{Potential inside } V(r, \theta, \varphi) = \sum_{l,m} A_{lm} r^l Y_{lm}(\theta, \varphi)$$

$$\text{at } r=a \quad V = V_0(\theta, \varphi) = \sum_{l,m} A_{lm} a^l Y_{lm}(\theta, \varphi)$$

$$\leadsto A_{lm} = \int \frac{V_0(\theta', \varphi')}{a^l} Y_{lm}^*(\theta', \varphi') d\Omega'$$

$$\leadsto V(r, \theta, \varphi) = \sum_{l,m} \left[\int \frac{V_0(\theta', \varphi')}{a^l} Y_{lm}^*(\theta', \varphi') \cdot r^l Y_{lm}(\theta, \varphi) \right] d\Omega'$$

$$= \sum_{l,m} V_0(\theta', \varphi') \left(\frac{r}{a}\right)^l \cdot Y_{lm}^*(\theta', \varphi') \cdot Y_{lm}(\theta, \varphi) d\Omega' \quad \leftarrow \textcircled{1}$$

for $r \rightarrow z$: $y=0$; $\theta \rightarrow 0$

$$\begin{aligned} \frac{a \cdot (a^2 - z^2)}{(z^2 + a^2 - 2az \cos \gamma)^{3/2}} &= \frac{a^3 (1 - z^2/a^2)}{a^3 \left[1 - \frac{z}{a} \right]^{3/2}} = \frac{1 + \frac{z}{a}}{\left(1 - \frac{z}{a}\right)^2} = \left(1 + \frac{z}{a}\right) \left(1 - \frac{z}{a}\right)^{-2} \\ &= \left(1 + \frac{z}{a}\right) \cdot \left[1 + 2\frac{z}{a} + 3\left(\frac{z}{a}\right)^2 + \dots \right] \\ &= 1 + 3 \cdot \frac{z}{a} + 5 \cdot \left(\frac{z}{a}\right)^2 + \dots \\ &= \sum_l (2l+1) \left(\frac{z}{a}\right)^l \end{aligned}$$

Now let $z \rightarrow r$; multiply by $P_l(\cos \gamma)$ and let $\gamma \rightarrow \gamma$

$$\frac{a(a^2 - r^2)}{(r^2 + a^2 - 2ar \cos \gamma)^{3/2}} = \sum_l (2l+1) \left(\frac{r}{a}\right)^l P_l(\cos \gamma)$$

$$\begin{aligned} P_l(\cos \gamma) &= \frac{4\pi}{2l+1} \sum_m Y_{lm}^*(\gamma', \varphi') Y_{lm}(\gamma, \varphi) \\ &= \sum_l \sum_m \frac{4\pi}{(2l+1)} \cdot (2l+1) \cdot \left(\frac{r}{a}\right)^l \cdot Y_{lm}^* Y_{lm} \\ &= \sum_{lm} 4\pi \left(\frac{r}{a}\right)^l Y_{lm}^*(\gamma', \varphi') Y_{lm}(\gamma, \varphi) \end{aligned}$$

$$\begin{aligned} \Leftrightarrow V_c(\vec{r}) &= \frac{1}{4\pi} \int V_0(\theta', \varphi') \sum_{l,m} 4\pi \left(\frac{r}{a}\right)^l \cdot Y_{lm}^* Y_{lm} d\Omega' \\ &= \sum_{l,m} \underbrace{\int V_0(\theta', \varphi') \left(\frac{r}{a}\right)^l Y_{lm}^* Y_{lm} d\Omega'}_{= R_{lm}} = \sum_{l,m} \left(\frac{r}{a}\right)^l Y_{lm} \underbrace{\int V_0(\theta', \varphi') Y_{lm}^* d\Omega'}_{= R_{lm}} \\ &= \sum_{l,m} R_{lm} \left(\frac{r}{a}\right)^l \cdot Y_{lm}(\gamma, \varphi) = \text{equal to } \textcircled{1} \end{aligned}$$