



## Solutions for HW # 3

1) Problem 1.4, Jackson textbook:

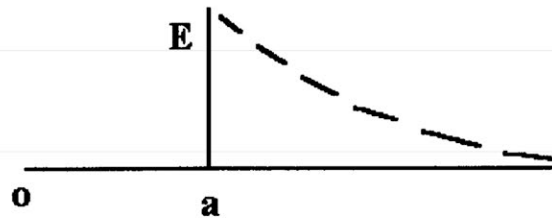
Gauss's Law:

$$\int \vec{E} \cdot d\vec{a} = \frac{Q_{\text{enclosed}}}{\epsilon_0}$$

a) Conducting sphere: all of the charge is on the surface  $\sigma = \frac{Q}{4\pi a^2}$ 

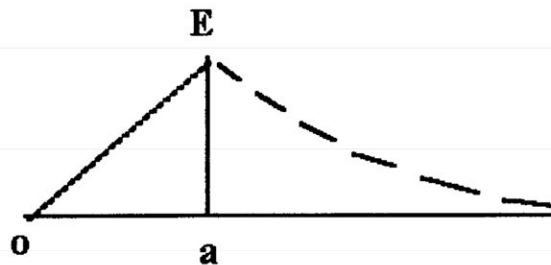
$$E4\pi r^2 = 0, \quad r < a \quad E4\pi r^2 = \frac{Q}{\epsilon_0}, \quad r > a$$

$$E = 0, \quad r < a \quad \vec{E} = \frac{Q}{4\pi\epsilon_0 r^2} \hat{r}, \quad r > a$$

b) Uniform charge density:  $\rho = \frac{Q}{\frac{4}{3}\pi a^3}$ ,  $r < a$ ,  $\rho = 0$ ,  $r > a$ .

$$E4\pi r^2 = \frac{Qr^3}{\epsilon_0 a^3} \rightarrow E = \frac{Qr}{4\pi\epsilon_0 a^3}, \quad r < a$$

$$E4\pi r^2 = \frac{Q}{\epsilon_0} \rightarrow E = \frac{Q}{4\pi\epsilon_0 r^2}, \quad r > a$$

c)  $\rho = A r^n$ 

$$Q = 4\pi \int r^2 dr A r^n = 4\pi A a^{n+3}/(n+3) \rightarrow \rho = \frac{(n+3)Q}{4\pi a^{n+3}} r^n, \quad r < a$$

$$\rho = 0, \quad r > a$$

$$E4\pi r^2 = \frac{(n+3)Q}{4\pi\epsilon_0 a^{n+3}} 4\pi r^{n+3}/(n+3) \rightarrow E = \frac{Q}{4\pi\epsilon_0 r^2} \left( \frac{r^{n+3}}{a^{n+3}} \right), \quad r < a$$

$$E = \frac{Q}{4\pi\epsilon_0 r^2}, \quad r > a$$

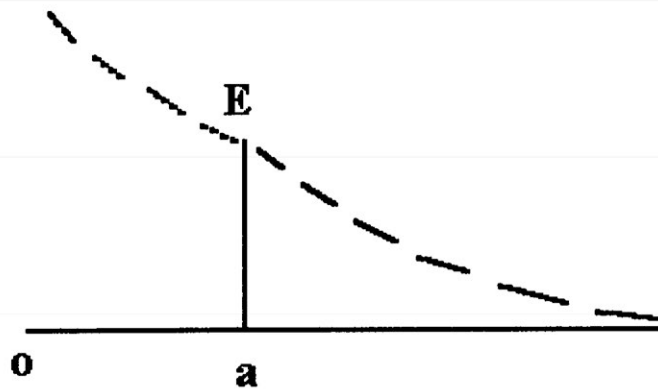
1.  $n = -2$ .

$$E = \frac{Q}{4\pi\epsilon_0 r a}, \quad r < a$$



1) Problem 1.4 - continued:

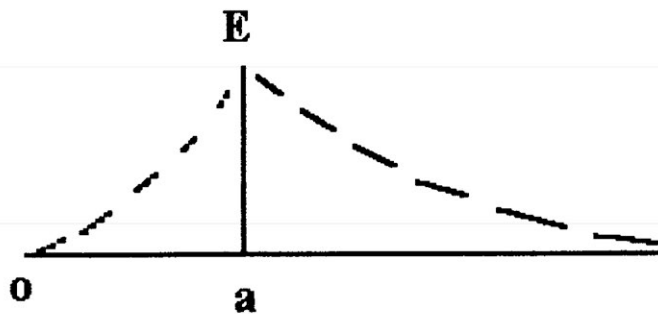
$$E = \frac{Q}{4\pi\epsilon_0 r^2}, \quad r > a$$



2.  $n = 2$ .

$$E = \frac{Qr^3}{4\pi\epsilon_0 a^5}, \quad r < a$$

$$E = \frac{Q}{4\pi\epsilon_0 r^2}, \quad r > a$$



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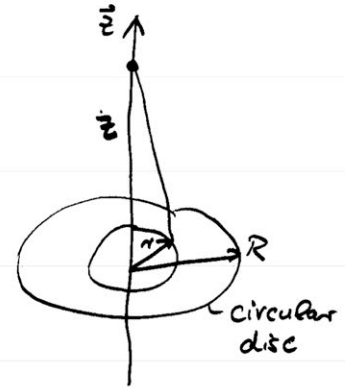
**Problem 2)** Find the potential and the electric field strength along the axis of a thin uniformly charged circular disc of radius  $R$  and total charge  $q$ . .....

$$1.) \quad \psi(z) = \frac{1}{4\pi\epsilon_0} \int_S \frac{d^2\vec{r}' \cdot \sigma(\vec{r}')}{|\vec{r} - \vec{r}'|}$$

$$= \frac{\sigma}{4\pi\epsilon_0} \int_S \frac{d^2\vec{r}'}{|\vec{r} - \vec{r}'|}$$

$$= \frac{\sigma}{4\pi\epsilon_0} \int_S \frac{d^2\vec{r}'}{\sqrt{z^2 + r'^2}} = \frac{\sigma}{4\pi\epsilon_0} \int_S \frac{2\pi r' dr'}{\sqrt{z^2 + r'^2}}$$

$$= \frac{\sigma}{2\epsilon_0} \int_0^R \frac{r' dr'}{\sqrt{z^2 + r'^2}} = \frac{\sigma}{2\epsilon_0} \sqrt{z^2 + r'^2} \Big|_0^R = \frac{\sigma}{2\epsilon_0} [\sqrt{z^2 + R^2} - \sqrt{z^2}]$$



$$\hookrightarrow \psi(z) = \begin{cases} \sigma/2\epsilon_0 \cdot (\sqrt{z^2 + R^2} - z) & \text{for } z > 0 \\ \sigma/2\epsilon_0 \cdot (\sqrt{z^2 + R^2} + z) & \text{for } z < 0 \end{cases}$$

2.) From symmetry:  $E_\phi = E_r = 0$  on axis

$$\hookrightarrow E_z = \frac{\partial \psi}{\partial z} = \begin{cases} \sigma/2\epsilon_0 (1 - \frac{z}{\sqrt{z^2 + R^2}}) & \text{for } z > 0 \\ \sigma/2\epsilon_0 (1 + \frac{z}{\sqrt{z^2 + R^2}}) & \text{for } z < 0 \end{cases}$$

$$3.) \text{ for } z \rightarrow +\infty \quad E_z \rightarrow \frac{\sigma}{2\epsilon_0} (1 - (1 + (\frac{R}{z})^2)^{-1/2}) \approx \frac{\sigma}{2\epsilon_0} (1 - 1 + \frac{1}{2} \frac{R^2}{z^2}) \\ = \frac{\sigma R^2}{4\epsilon_0 z^2} = \frac{q^2}{4\pi\epsilon_0 z^2} = \text{field of point charge}$$

$$\text{Similar } z \rightarrow -\infty \quad E_z \rightarrow -\frac{q^2}{4\pi\epsilon_0 z^2}$$

$$4.) E(z^+) - E(z^-) = \frac{\sigma}{2\epsilon_0} - (-\frac{\sigma}{2\epsilon_0}) = \frac{\sigma}{\epsilon_0}$$

$$5.) \text{ near disc: } E_z|_{z \rightarrow 0^+} = \left( \frac{\sigma}{2\epsilon_0} \cdot 1 - \frac{z}{\sqrt{z^2 + R^2}} \right) \Big|_{z \rightarrow 0} = \frac{\sigma}{2\epsilon_0} - \text{field of infinite plane.}$$

**Problem 3)** Consider a spherically symmetric charge distribution  $\rho = \rho(r)$ . By dividing the charge distribution into spherical shells, find the potential and the electric field ...

a.) Due to the spherical symmetry of the problem, we have

$$\vec{E} = E_r \cdot \hat{r}, \quad E_\theta = E_\phi = 0$$

$$E_r = E_r(r)$$



with Gaussian surface: - sphere of  $r'$

a1)  $r < r'$ : Gauss's law:  $\oint_S \vec{E} \cdot d\vec{A} = q/\epsilon_0$

$q = 0$  (charge is outside sphere)

$$\hookrightarrow \oint_S \vec{E} \cdot d\vec{A} = \oint_S E_r dA = E_r(r) \cdot 4\pi r'^2 = 0 \Rightarrow \underline{E_r(r) = 0}$$

a2)  $r > r'$ :  $\oint_S \vec{E} \cdot d\vec{A} = E_r(r) \cdot 4\pi r'^2 = q/\epsilon_0 = \frac{\sigma \cdot 4\pi r'^2}{\epsilon_0}$

$$\hookrightarrow E_r(r) = \frac{\sigma}{\epsilon_0} \left(\frac{r'}{r}\right)^2 \Rightarrow E_r(r) = \begin{cases} 0 & \text{for } r < r' \\ \frac{\sigma}{\epsilon_0} \cdot \left(\frac{r'}{r}\right)^2 & \text{for } r > r' \end{cases}$$

b.) potential:

$$\psi(r) = \psi(\infty) + \int_r^\infty E(r) dr$$

b1:  $r > r'$ :  $\psi(r) = \int_r^\infty \frac{\sigma}{\epsilon_0} \left(\frac{r'}{r}\right)^2 dr = \frac{\sigma}{\epsilon_0} (r')^2 \cdot \frac{1}{r} \Big|_r^\infty = \underline{\frac{\sigma}{\epsilon_0} \frac{(r')^2}{r}}$

b2:  $r < r'$ :  $\psi(r) = \int_r^\infty E(r) dr = \underbrace{\int_r^{r'} E(r) dr}_{=0} + \int_{r'}^\infty E(r) dr = \frac{\sigma}{\epsilon_0} \frac{(r')^2}{r'} = \underline{\frac{\sigma \cdot r'}{\epsilon_0}}$

$$\hookrightarrow \psi = \begin{cases} \frac{\sigma}{\epsilon_0} \frac{(r')^2}{r} & \text{for } r > r' \\ \frac{\sigma \cdot r'}{\epsilon_0} & \text{for } r < r' \end{cases}$$

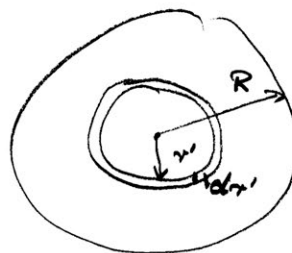
c.) Next, a spherically symmetric charge distribution  $\rho(r)$  of Radius  $R$  can be represented as a sum of infinitely thin spherical shells of radius  $r'$  ( $0 \leq r' \leq R$ ) and  $\Rightarrow$



Problem 3) continued ...

→ surface charge density

$$\sigma = \frac{4\pi (r')^2 d\tau' \rho(r')}{4\pi (r')^2} = \rho(r') d\tau' = \sigma(r')$$



Due to superposition principle, the total electric field and potential are calculated as integrals over  $r'$ .

1.)  $r > R$ :  $E_r(r) = \int_0^R dE_r(r')$ , where  $dE_r(r')$  is the field created by spherical shell of radius  $r'$  at point  $r$  and thickness  $dr'$ .

$$\hookrightarrow E_r(r) = \int_0^R \frac{\sigma(r')}{\epsilon} \left(\frac{r'}{r}\right)^2 dr' = \int_0^R \frac{1}{\epsilon} \rho(r') \frac{r'^2}{r^2} dr' = \frac{1}{\epsilon \cdot r^2} \int_0^R \rho(r') r'^2 dr'$$

$$\text{potential } \psi(r) = \int_0^R d\psi(r') = \int_0^R \frac{\sigma(r')}{\epsilon} \frac{(r')^2}{r} dr' = \int_0^R \frac{1}{\epsilon_0 r} \rho(r') (r')^2 dr'$$

2.)  $r < R$ :  $E_r(r) = \int_0^r dE_r(r') = \int_0^r \frac{1}{\epsilon} \rho(r') \left(\frac{r'}{r}\right)^2 dr' = \frac{1}{\epsilon r^2} \int_0^r \rho(r') r'^2 dr'$

and  $\psi(r) = \int_0^R d\psi(r') = \int_0^r d\psi(r') + \int_r^R d\psi(r')$

$$= \int_0^r \frac{\sigma(r') r'^2}{\epsilon \cdot r} + \int_r^R \frac{\sigma(r') r'}{\epsilon}$$

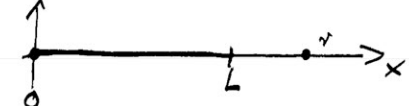
$$= \frac{1}{\epsilon_0 r} \int_0^r \rho(r') \cdot (r')^2 dr' + \frac{1}{\epsilon} \int_r^R \rho(r') r' dr'$$

$\hookrightarrow r > R$ :  $E_r(r) = \frac{1}{\epsilon r^2} \int_0^R \rho(r') r'^2 dr'$ ,  $\psi(r) = \frac{1}{\epsilon r} \int_0^R \rho(r') r'^2 dr'$

$r < R$ :  $E_r(r) = \frac{1}{\epsilon r^2} \int_0^r \rho(r') r'^2 dr'$ ,  $\psi(r) = \frac{1}{\epsilon r} \int_0^r \rho(r') r'^2 dr' + \frac{1}{\epsilon} \int_r^R \rho(r') r' dr'$

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**Problem 4)** A line conductor of length  $L$  and total charge  $Q$  lies on the  $x$ -axis with one end on the origin. Find the electric potential and the electric field ...

$$\rho(x) = \frac{Q}{L} \quad (0 \leq x \leq L)$$


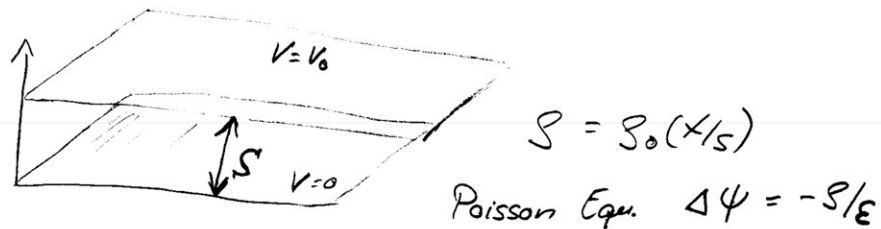
potential 
$$\psi(x) = \frac{1}{4\pi\epsilon} \int_0^L \frac{\rho(x') d^3x'}{|\vec{x}-\vec{x}'|} = \frac{1}{4\pi\epsilon} \frac{Q}{L} \int_0^L \frac{dx'}{|x-x'|}$$

$$= -\frac{1}{4\pi\epsilon} \frac{Q}{L} \ln(x-x') \Big|_0^L = -\frac{1}{4\pi\epsilon} \frac{Q}{L} \ln\left(\frac{x-L}{x}\right)$$

$$= -\frac{1}{4\pi\epsilon} \frac{Q}{L} \cdot \ln\left(1-\frac{L}{x}\right)$$

field 
$$\vec{E}(\vec{x}) = -\nabla\psi(x) = -\frac{1}{4\pi\epsilon} \frac{Q}{L} \left(\frac{1}{1-\frac{L}{x}}\right) \cdot \left(\frac{L}{x^2}\right) = \frac{1}{4\pi\epsilon} \cdot \frac{Q}{(x-L)x} \hat{x}$$

**Problem 5)** Two infinite parallel plates separated by a distance  $s$  are at the potentials 'zero' and  $V_0$ . a) find the potential ... and b) Find the surface charge densities on the plates.



a.) 
$$\frac{\partial^2\psi}{\partial x^2} = -\frac{\rho_0(x/s)}{\epsilon_0} ; \quad \sim \quad \frac{\partial\psi}{\partial x} = -\frac{\rho_0}{\epsilon_0} \cdot \frac{x^2}{2s} + C_1$$

$$\hookrightarrow \psi = -\frac{\rho_0}{\epsilon_0} \cdot \frac{x^3}{6s} + C_1 \cdot x + C_2$$

boundary:  $\psi(0) = 0 = C_2$

$$\psi(s) = V_0 = -\frac{\rho_0}{\epsilon_0} \cdot \frac{s^3}{6s} + C_1 \cdot s \quad \sim \quad C_1 = \frac{1}{s} \left( V_0 - \frac{\rho_0 \cdot s^2}{6} \right)$$

$$\hookrightarrow \psi(x) = \frac{-\rho_0}{\epsilon_0} \cdot \frac{x^3}{6s} + \left( \frac{V_0}{s} + \frac{\rho_0 \cdot s}{6\epsilon_0} \right) \cdot x$$

b) 
$$\vec{E} = -\nabla\psi ; \quad \vec{E} = -\frac{\partial}{\partial x} \left( \frac{-\rho_0}{\epsilon_0} \frac{x^3}{6s} + \left( \frac{V_0}{s} + \frac{\rho_0 \cdot s}{6\epsilon_0} \right) x \right)$$

$$= -\frac{\rho_0}{\epsilon_0} \frac{x^2}{2s} + \frac{V_0}{s} + \frac{\rho_0 \cdot s}{6\epsilon_0}$$

for  $x=0$ :  $E_x = \frac{\sigma}{\epsilon_0} \quad \sim \quad \sigma = \frac{-\rho_0 x^2}{2s} + \frac{V_0 \cdot \epsilon_0}{s} + \frac{\rho_0 \cdot s}{6}$ , for  $x=s$ :  $E_x = -\frac{\sigma}{\epsilon_0}$