

## Spectral Interdependency Methods

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### Synonyms

[Coherence and Granger causality spectral analysis](#); [Multivariate spectral analysis](#); [Oscillatory network activity analysis](#)

### Definition

Spectral interdependency methods are a means of statistically quantifying the interrelationship between a pair of dynamic processes as a function of frequency or time period of oscillation. The measures of spectral interdependency are derived from the time series recordings of dynamic systems either by using autoregressive modeling (parametric method) or by using direct Fourier or wavelet transforms (nonparametric method). For a pair of multivariate stationary processes (1 and 2), there are three measures that characterize the spectral interdependency between these processes. They are total interdependence ( $M_{1,2}$ ), Granger causality (one-way effect or directional influence from the first process to the second process,  $M_{1 \rightarrow 2}$ , or from the second to the first,  $M_{2 \rightarrow 1}$ ), and instantaneous causality (measure of reciprocity,  $M_{1,2}$ ). In general, the total interdependence is the sum of directional influences and instantaneous causality frequency by frequency ( $M_{1,2} = M_{1 \rightarrow 2} + M_{2 \rightarrow 1} + M_{1,2}$ ) and is related to coherence ( $C_{12}$ ) as  $M_{1,2} = -\ln(1 - C_{12})$ . These measures can also be estimated for nonstationary processes by using moving time-windowed autoregressive modeling or Fourier transformations or wavelet transformations. Coherence and spectral Granger causality are well-accepted measures in neuroscience to characterize frequency-specific interdependence between multiple time series from multisite neurophysiological recordings.

### Detailed Description

Many processes in nature, including brain processes, have oscillatory motion, and the time series measurements of their activity are rich in oscillatory content, lending them naturally to spectral analysis. Spectral interdependency methods are used to study the relationship in the frequency domain between multiple processes from their simultaneously recorded time series signals. Consider a pair of zero-mean, stationary processes (1 and 2) in which simultaneously measured time series at a sampling rate of  $f_s$  are represented as 1:  $X_1(1), X_1(2), \dots, X_1(t), \dots$  and 2:  $X_2(1), X_2(2), \dots, X_2(t), \dots$ . The spectral interdependency measures as defined above are derived from the spectral matrix (S) and/or from the transfer function (H) and noise covariance matrix ( $\Sigma$ ), which can be estimated by the parametric (Ding et al. 2006) or nonparametric approaches applied to these time series

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(Dhamala et al. 2008a, b). Brief mathematical derivations and descriptions of these interdependency measures ( $C_{12}$ ,  $M_{1,2}$ ,  $M_{1 \rightarrow 2}$ ,  $M_{2 \rightarrow 1}$  and  $M_{1,2}$ ) are included below.

Directed transfer function (DTF) (Kaminski et al. 2001) and partial directed coherence (PDC) (Baccala and Sameshima 2001) are the accepted alternative measures of directional influence, equivalent to the measures for  $M_{1 \rightarrow 2}$  and  $M_{2 \rightarrow 1}$  as defined above. DTF is obtained from H and PDC from the Fourier transform of model coefficients in the parametric approach.

**Parametric Approach.** Jointly,  $X_1$  and  $X_2$  series can be represented as the following bivariate autoregressive (AR) models:

$$\begin{aligned} X_1(t) &= \sum_{j=1}^{\infty} a_{11,j} X_1(t-j) + \sum_{j=1}^{\infty} a_{12,j} X_2(t-j) + \varepsilon(t) \\ X_2(t) &= \sum_{j=1}^{\infty} a_{21,j} X_1(t-j) + \sum_{j=1}^{\infty} a_{22,j} X_2(t-j) + \eta(t), \end{aligned} \quad (1)$$

where  $\varepsilon$  and  $\eta$  are residual (one-step-ahead prediction) errors and are uncorrelated over time. After Fourier transforming the bivariate AR representation (1) and applying proper ensemble average, we obtain the following spectral density matrix S as a function of frequency (f):

$$S(f) = H(f) \Sigma H^*(f), \quad (2)$$

where \* denotes the matrix adjoint. Here, the noise covariance matrix  $\Sigma$  is computed from the residual errors  $\varepsilon(t)$ ,  $\eta(t)$  and the transfer function matrix  $\mathbf{H}(f)$  is constructed from the matrix inverse of the Fourier transforms of the coefficients a's:

$$\begin{aligned} \Sigma &= \begin{pmatrix} \text{var}(\varepsilon) & \text{cov}(\varepsilon, \eta) \\ \text{cov}(\varepsilon, \eta) & \text{var}(\eta) \end{pmatrix} \\ H_{lm}(f) &= \left( \delta_{lm} - \sum_{k=1}^{\infty} a_{lm,k} e^{-i2\pi f k} \right)^{-1}, \end{aligned} \quad (3)$$

where  $\delta_{lm}$  is the Kronecker delta function with the matrix element index  $lm$ .

**Nonparametric Approach.** S, H, and  $\Sigma$  can also be estimated by using the nonparametric spectral methods (Dhamala et al. 2008a, b) without explicitly fitting the time series  $X_1(t)$  and  $X_2(t)$  in autoregressive models. In this approach, S is constructed by Fourier transforming  $X_1$  and  $X_2$  and properly averaging over ensembles usually with multitapers (Mitra and Pesaran 1999):

$$S_{lm} = \langle x_l(f) x_m(f)^* \rangle, \quad (4)$$

where  $l_m$  is the index for time series and matrix element,  $\langle \rangle$  represents averaging over ensemble, and X's are the direct Fourier transforms of X's. H and  $\Sigma$  can be derived from the minimum-phase factors of S:

$$\begin{aligned} H &= \psi A_0^{-1} \\ \Sigma &= A_0 A_0^T, \end{aligned} \quad (5)$$

where T stands for matrix transposition,  $\psi(e^{i\theta}) = \sum_{k=0}^{\infty} A_k e^{ik\theta}$  defined on the unit circle and

$$A_k = \frac{1}{2\pi} \int_{-\pi}^{\pi} \psi(e^{i\theta}) e^{-ik\theta} d\theta.$$

**Spectral Interdependence.** The coherence function  $C(f)$  is derived from the cross spectra normalized by the product of the individual auto spectra and measures the amount of interdependence (synchrony) as a function of frequency:

$$C(f) = \frac{|S_{12}(f)|^2}{S_{11}(f)S_{22}(f)}. \quad (6)$$

$C(f)$  is sensitive to both amplitude and phase relationships between processes at  $f$  and its value ranges from 0 (no interdependence) to 1 (maximum interdependence).  $C(f)$  is related to Geweke's measure of total interdependence ( $M_{1,2}$ ) (Ding et al. 2006):

$$M_{1,2}(f) = -\ln(1 - C(f)), \quad (7)$$

whose value ranges from 0 to infinity. The total spectral interdependence ( $M_{1,2}$ ) is equal to the sum of Granger causality or directional influences (one-way effects,  $M_{1 \rightarrow 2}$  and  $M_{2 \rightarrow 1}$ ) and instantaneous causality ( $M_{1,2}$ ) (Geweke 1982):

$$M_{1,2} = M_{1 \rightarrow 2} + M_{2 \rightarrow 1} + M_{1,2} \quad (8)$$

Here, directional influences between 1 and 2 are given by (Geweke 1982)

$$M_{1 \rightarrow 2}(f) = \ln \frac{S_{22}(f)}{\tilde{H}_{11}(f) \sum_{11} \tilde{H}_{11}^*(f)} \quad M_{2 \rightarrow 1}(f) = \ln \frac{S_{11}(f)}{\tilde{H}_{22}(f) \sum_{22} \tilde{H}_{22}^*(f)}, \quad (9)$$

where  $\tilde{H}_{11} = H_{11} + \sum_{11}^{12} H_{12}$ ,  $\tilde{H}_{22} = H_{22} + \sum_{22}^{12} H_{21}$ , and instantaneous causality  $M_{1,2}$  is given by

$$M_{1,2}(f) = -\ln \frac{|S(f)|}{(\tilde{H}_{11}(f) \sum_{11} \tilde{H}_{11}^*(f)) (\tilde{H}_{22}(f) \sum_{22} \tilde{H}_{22}^*(f))}. \quad (10)$$

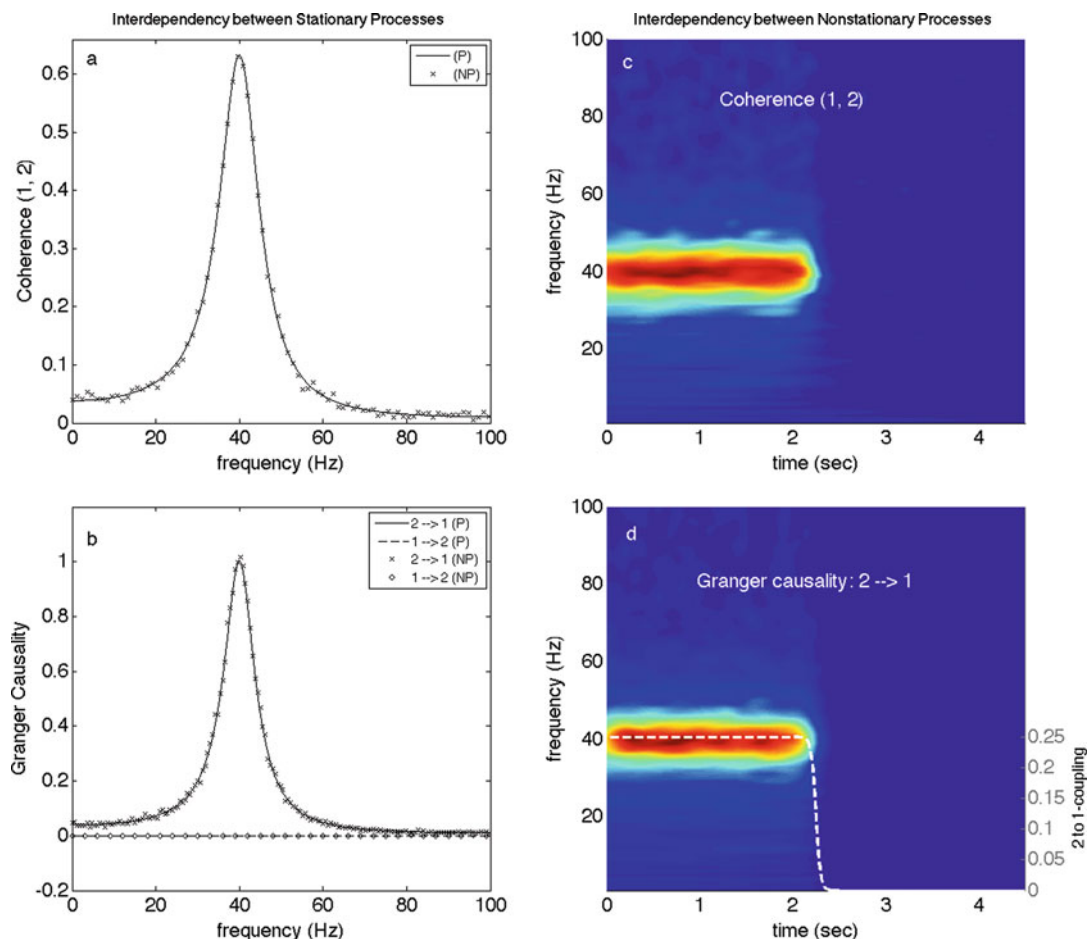
Granger causality at low frequencies between a pair of cointegrated series depends only on a few statistically interpretable coefficients from the error correction model if Hosoya's spectral decomposition (Hosoya 1991) is used (Granger and Lin 1995).

The time-domain counterparts of these spectral measures in Eq. 8 are obtained by integrating the spectral measures over the entire frequency range as  $\frac{2}{f_s} \int_0^{\frac{f_s}{2}} M(f) df$ . The integral of total interdependence measures the total amount of mutual information, and the other integrals are simply the respective time-domain measures of Granger causality.

**Spectral Interdependence for Nonstationary Processes.** The spectral interdependency measures can also be defined for nonstationary processes by treating their time series in sufficiently short windows as locally stationary processes. They can be estimated as a function of time and frequency by using moving time-windowed autoregressive modeling (Liang et al. 2000), moving time-windowed Fourier transformations (Shiogai et al. 2012), or wavelet transformations (Dhamala et al. 2008a, b).

**Numerical Examples.** Here, we generate time series (500 trials) from jointly stationary and nonstationary processes and illustrate the estimation of coherence and Granger causality spectra using the parametric (P) and nonparametric (NP) approaches. We consider two interacting autoregressive processes, 2:  $\{X_2(t)\}$  driving 1:  $\{X_1(t)\}$ , similar to the network model considered in Dhamala et al. 2008b, where  $X_1 = 0.55 X_1(t - 1) - 0.8X_1(t - 2) + C(t)X_2(t - 1) + \varepsilon(t)$  and  $X_2(t) = 0.55X_2(t - 1) - 0.8X_2(t - 2) + \eta(t)$ . Here,  $(\varepsilon(t), \eta(t))$ 's are independent white noise processes with zero means and unit variances, the sampling rate is considered to be 200 Hz, and the coupling strength  $C(t)$  remains 0.25 for time  $t$  in the stationary condition and slowly changes from 0.25 to 0 around  $t = 2$  s in the case of nonstationary processes. P and NP approaches agree well in coherence and Granger causality spectra (Fig. 1a, b). Morlet wavelet transform-based method yields complete time-frequency maps of coherence (Fig. 1c) and Granger causality (Fig. 1d), consistent with the trend of the 2 to 1 coupling as shown on the right side of Fig. 1d.

**Extensions of Spectral Interdependency Methods.** As an extension of the ordinary coherence described above, block coherence (Nedungadi et al. 2011) can estimate coherence spectra between pairs of nonoverlapping time series. Conditional Granger causality (Geweke 1984; Hosoya 2001) measures directional influences between two processes eliminating the effect of a third process,



**Fig. 1** Coherence (a, c) and Granger causality (b, d) spectra between stationary (a, b) and nonstationary (c, d) processes (1 and 2). Parametric (P) and Fourier transform-based nonparametric (NP) methods are used here to evaluate these quantities (a, b) between the stationary processes. The wavelet transform-based time-frequency maps of coherence (c) and Granger causality (d) recover the time-varying nature of 2 to 1 coupling (shown on the right y-axis in d)

thereby distinguishing between direct and mediated causality (Ding et al. 2006; Dhamala 2008b). The multivariate version of the spectral Granger causality between two processes makes use of this idea of eliminating the causal effect from all other interrelated processes as an extension of the conditional Granger causality defined for three processes.

**Applications in Neuroscience.** Spectral interdependency measures have been instrumental in attempts to understand the relationships between oscillatory brain processes at various spatial scales from multisite brain activity recordings. Coherence is widely used in neuroscience (Siegel et al. 2012; Roberts et al. 2013). The insights provided by this measure have even led to the “neuronal communication through coherence” hypothesis (Fries 2005; Roberts et al. 2013). Spectral Granger causality and equivalent directional measures have been used in a variety of brain signal recordings, such as local field potentials, EEG, MEG, and fMRI, in animals and humans (see Bressler and Seth 2011; Friston et al. 2012 for reviews). Because of the unknown theoretical distributions of coherence and spectral Granger causality, establishing statistical significance in these measures derived from experimental time series requires data resampling (surrogate) methods such as jackknifing (Bokil et al. 2010), bootstrapping, and random permutation (Seth 2010).

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